

Molecular dynamics simulation on temperature and stain rate-dependent tensile response and failure behavior of Ni-coated CNT/Mg composites

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Abstract

This paper focuses on the molecular dynamics (MD) simulation of the tensile response of Ni-coated CNT-reinforced magnesium matrix composites (Ni-CNT/Mg) subjected to uniaxial tension at different temperatures and strain rates. The results show that Ni-CNTs can improve the mechanical properties of the composites effectively. The maximum stress of Ni-(6,6) CNT/Mg is, respectively, increased by 25.66 and 11.13%, while the elastic modulus is increased by 23.69 and 14.43% compared with those of the single-crystal Mg and uncoated $(6,6)$ CNT/Mg at 300 K and 1×10^9 s⁻¹. In addition, the calculated elastic modulus of the Ni-(6,6)CNT/Mg composite is consistent with the prediction based on the rule-of-mixture. The Ni-CNT/Mg composites still have better mechanical properties at 500 K but exhibit a significant temperature softening effect in the temperature range 100–500 K and a strong positive strain-rate sensitivity at a strain rate greater than 1×10^9 s⁻¹. The various failure modes of the composite at the nano-scale are mainly determined by the combined effects of the different factors such as the atomic disordering, the void and dislocation nucleation, structural phase transformations in Mg nanocrystals near the interface, and the subsequent fracture of the Ni-coated CNT close to the voids.

1 Introduction

Novel lightweight magnesium (Mg) alloys and Mg matrix nanocomposites are now considered having broad application prospects in the automobile, aerospace, and weapon industries due to their high strength-to-weight ratio and better combination of mechanical properties and energy absorption characteristics [[1](#page-10-0)]. Currently, the research on

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mechanical properties and failure mechanisms of nanosized particulate and whisker reinforced Mg-based composites at low strain rates has achieved some results [\[2](#page-10-1)]. Compared with nanoparticles and nanofibers, carbon nanotubes (CNTs) have been expected as ideal reinforcements for composites [\[3](#page-10-2)] due to their high aspect ratios, high strength, high toughness, and lightweight characteristics. As a result, CNT-reinforced Mg matrix composites possess the potential to be the next generation structural materials, so it is necessary to quantificationally analyze the bonding characteristics between CNTs and Mg matrix and their influences on the macroscopic dynamic mechanical properties.

The uniform CNT dispersion and strong interface bonding are two key issues to obtain high strength and good ductility of Mg matrix composites reinforced by CNTs, while the interfacial bonding is one of the most significant challenging issues that determines the efficiency of load transfer from the matrix to the CNTs. Existing experimental studies on the CNT/Mg interface are mostly focused on the interface microstructure characterization and interface chemistry as well as their qualitative effects on the macroscopic mechanical properties of the composites [[4](#page-10-3), [5](#page-10-4)]. Liang et al. [[4\]](#page-10-3) fabricated 0.5 and 1.0 wt% Ni-CNT-reinforced AZ91D Mg alloy (Ni-CNTs/AZ91D) composites by ultrasonic

vibration method followed by extrusion. They found that the ultimate tensile strength (UTS) of the composite was increased by 17% after adding 0.5 wt% CNTs, while the UTS of the 1.0 wt% CNTs/AZ91D composite was less than that of AZ91D alloy. They believe that the fracture mode is the interface disengaged mechanism, but the interface between Ni-CNTs and Mg alloy requires further study. Investigations by Nai et al. [[5\]](#page-10-4) found that the presence of nickel coating on the CNT results in the formation of Mg_2Ni intermetallics at the Ni-CNT/Mg interface, enhances the grain refinement and uniform dispersion of the Ni-CNT reinforcements in the Mg matrix, and so improves the micro-hardness, UTS and 0.2% yield strength of the monolithic Mg by 41, 39, and 64%, respectively.

Molecular dynamics (MD) simulation technique has been widely used in the study of interface behavior, because accurate interface measuring methods are difficult to be carried out and the interface description theory is not yet complete. Zhou et al. [[6](#page-10-5)] studied the static tensile mechanical properties of Ni-CNT/Mg composite using MD simulation method. Their simulation results have verified and confirmed the effectiveness of Ni-CNTs in improving the mechanical properties of CNT/Mg composites. Song and Zha [\[7](#page-10-6)] analyzed the uniaxial tension behavior of Ni-CNT-reinforced gold matrix composite using MD simulation method. They found that the mechanical properties of Ni-CNT/Au composites were obviously improved compared with those of the uncoated CNT-reinforced composite and they believed that nickel coating could provide an effective channel for interfacial stress transfer. Duan et al. [[8\]](#page-10-7) investigated the damping characteristics of Ni-CNT-reinforced copper-matrix nanocomposites at the atomic level and demonstrated that the distorted Cu lattice structure caused by the attraction of Ni is the dominant factor for the high damping rate of Ni-coated CNT/Cu composites.

In the present paper, to study improvement mechanism in mechanical properties and damage evolution behavior of Nicoated CNT-reinforced Mg matrix composites undergoing dynamic loading conditions, MD simulation was conducted to investigate the tensile properties and failure behavior of Ni-coated armchair single-walled CNT/Mg composites under uniaxial tension by considering the coupled effects of temperature and strain rate. For convenience, the armchair SWCNT mentioned in this paper is denoted as CNT.

2 MD simulation approach

2.1 Simulation models

MD simulations of uniaxial tension of Ni-CNT/Mg composites at different temperatures and strain rates were carried out using the Large-scale Atomic/Molecular Massively Parallel Simulator (LAMMPS) software [[9\]](#page-10-8) in this work. The MD simulation model of the Ni-(6,6)CNT/Mg composite is shown in Fig. [1](#page-1-0). To represent the nano-scale structure inside a Ni-CNT/Mg composite, periodic boundary conditions (PBCs) were applied in the *x, y*, and *z* directions, while the *x, y*, and *z* axes are along the crystal orientation of [2*̄*110], [112*̄*0] and [0001] in hexagonal close packed (HCP) Mg single crystals,

Fig. 1 Atomic model of Ni-(6,6)CNT/Mg composite with **a** top view, **b** side view, and **c** embedded Ni-CNT in Mg. The red, blue, and yellow atoms denote magnesium, carbon, and nickel, respectively

Table 1 Geometrical parameters of armchair CNTs embedded in Mg matrix

Chirality index	(4,4)	(6.6)	(8,8)	(10,10)
Diameter (nm)	0.542	0.814	1.805	1.356
Length (nm)	9.8	9.8	9.8	9.8
Atom numbers	640	960	1280	1600

respectively. In addition, the initial equilibrium positions of Ni atoms on CNT surfaces are all near the centers of carbon-atom hexagonal rings in CNTs [\[10](#page-10-9)], and the distance between Ni atoms and CNT surfaces is about 0.15 nm.

All composite models have the same Mg matrix geometry with various inner diameters of Ni-coated armchair CNTs, and the composite models have equal dimensions: length of 9.80 nm, width of 6.67 nm, and depth of 6.44 nm. The geometrical parameters of CNTs used in the MD models are shown in Table [1](#page-2-0) and the CNTs with or without surface-coated nickel atoms are embedded in the center of the Mg matrix parallel to the axial tensile direction. Each model for Ni-coated (4,4) CNT, (6,6)CNT, (8,8)CNT, and (10,10)CNT-reinforced Mg composites contains 17,966 atoms (C: 640, Ni: 322, and Mg: 17,004), 18,160 atoms (C: 960, Ni: 480, and Mg: 16,720), 18,356 atoms (C: 1280, Ni: 640, and Mg: 16,436), and 18,550 atoms (C: 1600, Ni: 800, and Mg: 16,150), respectively.

2.2 Inter‑atomic interaction potentials

For MD simulation, the selection of potential function is a key factor to obtain reasonable results. In the present MD model, different potential functions were employed to describe the interactions of atoms. Herein, the empirical embedded-atom method (EAM) potential developed by Sun et al. [\[11](#page-10-10)] was adopted to describe the interactions between Mg atoms. The total energy E_i of an atom *i* is given by

$$
E_i = F_\alpha \left(\sum_{j \neq i} \rho_\beta(r_{ij}) \right) + \frac{1}{2} \sum_{j \neq i} \phi_{\alpha\beta}(r_{ij}), \tag{1}
$$

where F_a is the embedding energy, which is a function of the atomic electron density ρ_{β} ; $\phi_{\alpha\beta}$ is a pair potential interaction, and *α* and *β* are the element types of atoms *i* and *j*. Both summations in the formula are over all neighbors *j* of atom *i* within the cut-off distance.

The C–C inter-atomic interactions were simulated using the adaptive intermolecular reactive bond-order (AIREBO) [\[12\]](#page-10-11) potential with a cut-off radius of 2.0 Å [[13\]](#page-10-12). The AIREBO potential is composed of the Reactive Empirical Bond-Order (REBO) potential, the Lennard–Jones (LJ) potential and the torsional interaction potential. It is expressed as

$$
E = \frac{1}{2} \sum_{i} \sum_{j \neq i} \left[E_{ij}^{\text{REBO}} + E_{ij}^{\text{LJ}} + \sum_{k \neq i,j} \sum_{l \neq i,j,k} E_{kijl}^{\text{TORSION}} \right], \qquad (2)
$$

where the E_{ij}^{REBO} item describes short-ranged C–C interactions with distance $r < 2.0$ Å. The E_{ij}^{LJ} item adds longer ranged interactions (2.0 $\AA < r <$ cut-off distance) using the standard 12-6 LJ potential. The E_{ij}^{TORSION} term is an explicit four-body potential that describes various dihedral angle preferences in hydrocarbon configurations.

The C–Ni and Ni–Ni interaction potential parameters were determined using the bond-order potential [\[14\]](#page-10-13). The functions of the bond-order potential are shown as follows:

$$
E_{ij}^{C-Ni} = \frac{D_e}{S-1} \exp \left\{-\beta \sqrt{2S} (r_{ij} - R_e) \right\}
$$

- $B^* \frac{D_e S}{S-1} \exp \left\{-\beta \sqrt{2/S} (r_{ij} - R_e) \right\}$ (3)

$$
B^* = \left\{1 + b\left(N^{\rm C} - 1\right)\right\}^{\delta} \tag{4}
$$

$$
E_{ij}^{\text{Ni-Ni}} = \frac{D_e}{S-1} \exp\left\{-\beta \sqrt{2S} (r_{ij} - R_e)\right\}
$$

$$
-\frac{D_e S}{S-1} \exp\left\{-\beta \sqrt{2/S} (r_{ij} - R_e)\right\}
$$
(5)

$$
D_e = D_{e1} + D_{e2} \exp\{-C_D(N^M - 1)\}\tag{6}
$$

$$
R_e = R_{e1} - R_{e2} \exp\{-C_R (N^M - 1)\}\tag{7}
$$

$$
f(r_{ij}) = \begin{cases} 1 & r_{ij} < R_1 \\ \frac{1}{2} \left[1 + \cos \frac{\pi (r_{ij} - R_1)}{R_2 - R_1} \right] & R_1 < r_{ij} < R_2 \\ 0 & r_{ij} > R_2 \end{cases}
$$
(8)

$$
N_i^{\rm C} = 1 + \sum_{\text{carbon}} f(r_{ik})
$$
\n(9)

$$
N_i^{\mathbf{M}} = 1 + \sum_{\text{metal}} f(r_{ik}), \ N_{ij}^{\mathbf{M}} = \left(N_i^{\mathbf{M}} + N_j^{\mathbf{M}} \right) / 2, \tag{10}
$$

where E_{ij}^{C-Ni} and E_{ij}^{Ni-Ni} , respectively, denote the total potential energy in C–Ni and Ni–Ni atoms, and r_{ii} represents the distance between atoms *i* and *j*. The potential parameters *S,* β , *b*, and δ determine the shape of the potential function. The additional term B^* is a function of carbon coordination number N^C of a metal atom. The binding energy D_e and the equilibrium bond length *Re* are expressed as direct functions of the metal coordination number N^M . From the cut-off function $f(r_{ij})$, the coordination numbers of N_i^C and N_i^M can be obtained. The parameters used in this potential for the C–Ni and Ni–Ni interactions can be found in reference [[14\]](#page-10-13).

In addition, the interactions between Mg and C atoms or Mg and Ni atoms are modeled by the LJ potential $[6]$ $[6]$:

$$
E_{\text{LJ}}(r) = 4\varepsilon \left[\left(\frac{\sigma}{r} \right)^{12} - \left(\frac{\sigma}{r} \right)^6 \right],\tag{11}
$$

where ε is the depth of the potential well, σ is the distance where the inter-atomic potential is zero, and *r* is the distance between two atoms. The LJ (12-6) parameters are shown in Table [2](#page-3-0).

2.3 Simulation details and calculation of stress at nano‑scale

Before starting the MD simulations of uniaxial tension, the initial configurations must be energy minimized to obtain a stable structure by the conjugate gradient method. Thus, the systems were thermally equilibrated to the set temperature such as 300 K under the NPT ensemble with a pressure of 0 bar by running 40,000 steps with a time step of 1 fs. During the relaxation process, the temperature was kept constant at a certain temperature with Nose–Hoover thermostat and the system press was adjusted using Parrinello–Rahman barostat. Relaxation time of 40 ps was found to be enough for different Ni-CNT/Mg composite systems to reach the equilibrium stage, while the system temperature and total energy were found constant after 10 ps. Figure [2](#page-3-1) shows the energy convergence of four Ni-CNT/Mg composite systems with different CNT diameters at 300 K. After relaxation, these Ni-CNT/Mg composites were subjected to a uniaxial tensile loading. The tensile simulations were carried out with a pressure of 0 bar at different temperatures and strain rates in the positive direction of the *z*-axis (Fig. [1\)](#page-1-0).

During uniaxial tensile simulation, the velocity-Verlet algorithm was used for the integral calculation of the classical Newton's moving equation. The open-source visualization tool OVITO was utilized to visualize MD data and generate snapshots [[15](#page-10-14)]. The structure was characterized with the common neighbor analysis (CNA) technique to distinguish defects from HCP, FCC, and BCC environments.

In addition, the normal tensile stress at nano-scale has been calculated based on virial stress by the following equation:

$$
\sigma(\mathbf{r}) = \frac{1}{\Omega} \sum_{i} \left[-m_i \mathbf{v}_i \otimes \mathbf{v}_i + \frac{1}{2} \sum_{i \neq j} \mathbf{r}_{ij} \otimes \mathbf{f}_{ij} \right],
$$
 (12)

Fig. 2 Total energy variation during relaxation

where Ω is the total volume, and m_i and ν_i are the mass and velocity of atom *i*, respectively. ⊗ represents the cross product operator of two vectors and r_{ii} is the relative displacement vector from atom *i* to atom *j*, while f_{ii} is the force that the atom j acts on the atom i [[16\]](#page-10-15).

3 Results and discussion

3.1 Ni‑CNT effect on tensile properties of Mg and CNT/Mg composites

The tensile stress–strain curves of Ni-coated (6,6)CNT and uncoated (6,6)CNT-reinforced Mg composites as well as monolithic Mg at 300 K and 1×10^9 s⁻¹ are displayed in Fig. [3](#page-3-2) for comparison. In the case of the monolithic Mg,

Fig. 3 Stress–strain curves for Ni-(6,6)CNT/Mg,(6,6)CNT/Mg and monolithic Mg at 300 K and 1×10^9 s⁻¹

similar to the Ni-CNT/Mg composite model (Fig. [1b](#page-1-0)), the size of the monolithic Mg model was, respectively, $9.80\times6.67\times6.44$ nm in the *x*, *y*, and *z* directions, and PBCs were also assigned to the *x, y*, and *z* directions. It can be seen that the stress–strain curves for the above three materials have different characteristics. During the initial elastic stage and the subsequent nonlinear elastic stage, stress–strain curves are not overlapped, indicating that the elastic properties are sensitive to different materials. With the increase of applied strain, the Young's modulus decreases slightly, which can be seen from the slight change in slope towards the end of elastic deformation. The Young's modulus can be determined by linear regression analysis for the initial elastic stage [\[17\]](#page-10-16). Here, the Young's modulus for the Ni-CNT/Mg composites is about 83.12 GPa, while those for the CNT/Mg composites and the monolithic Mg are 77.01 and 67.20 GPa, respectively. It can be found that good agreement in elastic modulus of single-crystal Mg between the present MD prediction (67.20 GPa) and existing calculation result (69.6 GPa) in reference $[11]$ $[11]$ was achieved, preliminarily validating the applicability of our MD model. The enhancement in the elastic modulus of the Ni-CNT/Mg composites is directly correlated with the improved interface bonding of the Ni-CNT to the Mg matrix.

As mentioned before, the stress–strain curves at the initial elastic stages for the three materials (Fig. [3](#page-3-2)) differ a lot due to the contribution of the Ni-CNT to the elastic modulus. In fact, it can be easily seen that there is also a significant difference for the stress–strain curves after the stress reaches the highest point (at this point, the stress is defined as peak stress). For the Ni-CNT/Mg, the CNT/Mg composites, and the monolithic Mg, there all exist an ever-decreasing stage of the stress due to the nucleation of defects and a stress flow stage because of the interaction of defects after the stress reaches the peak stress, but a second large stress decrease stage occurs in the stress–strain curves for the first two due to the CNT fracture (at this point, the stress is defined as fracture stress) and the fracture propagation. Compared with the uncoated (6,6)CNT/Mg composite and the monolithic Mg matrix, the peak stress for the Ni-(6,6)CNT/Mg composite is, respectively, increased by 11.13 and 25.66%, while its fracture stress is separately increased by 28.47 and 223.64%, showing that the Ni-CNT also contributes significantly to the strength of the uncoated CNT/Mg composite and the Mg matrix.

To further investigate the effects of the inner diameter of CNT on the tensile behavior of Ni-CNTs and Ni-CNT/ Mg composites, Ni-coated armchair CNTs and Ni-CNT/ Mg matrix composites with different CNT chiral indexes are simulated, as shown in Figs. [4](#page-4-0) and [5,](#page-4-1) respectively. It should be noted that the tensile simulation of the Ni-CNT along its axial direction was conducted under the identical conditions with those of Ni-CNT/Mg composites including

Fig. 4 Stress–strain curves of Ni-CNTs with chirality of (4, 4), (6, 6), $(8, 8)$, and $(10, 10)$ at 300 K and 1×10^9 s⁻¹

Fig. 5 Stress–strain curves of Ni-CNT/Mg composites with different chiral indexes

temperature, strain rate, and Ni-CNT volumes except for boundary conditions. With respect to the Ni-CNT, an axial periodic boundary and free boundary conditions in the other two directions were adopted, since the mechanical properties of CNTs were insensitive to the simulation box under PBCs [18]. The stress–strain curves of $(4, 4)$, $(6, 6)$, $(8, 8)$, and $(10, 6)$ 10) Ni-CNTs are shown in Fig. [4](#page-4-0) in which the table lists the elastic moduli of the Ni-CNTs calculated by linear regression. It can be found from Fig. [4](#page-4-0) that the predicted elastic modulus by MD simulations depends on the CNT diameter. The elastic modulus of Ni-(6,6)CNT with the diameter of 0.814 nm predicted by MD simulation is 788.18 GPa and it is close to that reported in reference [[19](#page-10-18)] (747 GPa), thus further validating the present MD model. In addition, as can be observed from Fig. [5](#page-4-1), the mechanical behaviors such as the elastic modulus, the peak stress, and the fracture stress are significantly improved with the increase in the inner diameter of the Ni-CNT. However, the fracture strain increases with the increase in the inner diameter of the Ni-CNT. This is because the tensile stress in the nanotube is diameter dependent [\[18\]](#page-10-17) and the strengthening of the Ni-CNT/Mg composite is dominated by the Ni-CNTs reinforcement, while the fracture behavior of the Ni-CNT/ Mg composite maybe governed by more factors such as Ni-CNT size characteristics and deformation mechanisms [\[19](#page-10-18)].

It is worth noting that the elastic behavior of the Ni-CNT/ Mg composites is CNT diameter dependent and the elastic modulus of the Ni-CNT/Mg composites is considerably enhanced with the increase in the inner diameter of the Ni-CNT. However, because the mechanical behaviors of the Ni-CNT/Mg composites are not sensitive to the dimension of the simulation model by exertion of periodic boundaries, the effective elastic modulus of the Ni-CNT/Mg composites could still be estimated by a simple and well established rule-of-mixtures (ROM) [\[20](#page-10-19)] expressed as follows:

$$
E_{\text{Ni}-\text{CNT/Mg}} = E_{\text{Mg}} V_{\text{Mg}} + E_{\text{Ni}-\text{CNT}} V_{\text{Ni}-\text{CNT}},\tag{13}
$$

where $V_{\text{Ni-CNT}}$ and V_{Mg} are separately the volume fractions of the Ni-CNT and Mg matrix; $E_{\text{Ni-CNT/Mg}}$, $E_{\text{Ni-CNT}}$, and E_{Mg} are, respectively, the elastic moduli of the Ni-CNT/Mg composites, Ni-CNT, and Mg matrix. Here, the volume of the Ni-CNT can be approximated to be that of the CNT [\[21](#page-10-20)]. As such, the volume fractions of the Ni-CNT *V*_{Ni−CNT} and Mg matrix V_{Mg} in Ni-CNT/Mg composites can be, respectively,
c a l c u l a t e d as $V_{\text{Ni}-\text{CNT}} = \frac{2\pi R_{\text{CNT}}t \cdot l}{v - \left[\pi (R_{\text{CNT}} - t/2)^2 \cdot l\right]}$ a n d

 $V_{\text{Mg}} = 1 - V_{\text{Ni}-\text{CNT}}$, where R_{CNT} is the mean radius of the CNT, *t* is the wall thickness of the CNT and is about 0.34 nm, *l* is the length of the CNT, and *V* is the total volume of the Ni-CNT/Mg composite model.

Based on the MD elastic modulus predictions of a pure Mg matrix and an isolated Ni-coated SWCNT, the effective elastic moduli of the Ni-CNT/Mg composites with different Ni-CNT inner diameters corresponding to chirality of (4, 4), (6, 6), (8, 8), and (10, 10) CNTs are estimated by the rule-of-mixture and are illustrated in Fig. [6.](#page-5-0) Meanwhile, the ROM-based calculated moduli are compared with the MD modulus predictions of the Ni-CNT/Mg composites with varied CNT diameter sizes. It is found that there is good agreement of the elastic moduli between the ROM-based calculated results and the present MD predictions for different inner diameter of Ni-CNT-reinforced Mg composites. In addition, the present MD predictions also agree well with the experimental results of Liao and Yuan [\[22\]](#page-10-21).

3.2 Temperature effect

The temperature effect on tensile properties of Ni-(6,6) CNT/Mg composites at a strain rate of 1×10^9 s⁻¹ has been

Fig. 6 Comparison of elastic moduli predicted by MD and ROM for Ni-CNT/Mg composites with different Ni-CNT inner diameters corresponding to chirality of $(4, 4)$, $(6, 6)$, $(8, 8)$, and $(10,10)$ CNTs at 300 K and 1×10^9 s⁻¹

Fig. 7 Stress–strain curves of Ni-CNT/Mg composites at different temperatures and their comparisons with Mg single crystal and CNT/ Mg composite

investigated and the results are shown in Fig. [7](#page-5-1). It can be seen from Fig. [7](#page-5-1) that the tensile properties of Ni-CNT/Mg composite have a strong temperature dependence. When the temperature is increased from 100 to 500 K, the oscillations in simulated stress–strain curves for the Ni-CNT/Mg composite become larger and more irregular; however, the elastic modulus, the peak stress and the corresponding strain, and the fracture strain of the Ni-CNT/Mg composite both appear to decrease except for the fracture stress. Compared with the CNT/Mg composite and the monolithic Mg matrix, the Ni-CNT/Mg composite still shows superior mechanical properties at the high temperature of 500 K.

Figure [8](#page-6-0) shows the changes in tensile failure modes for Ni-CNT/Mg composites at the fracture strains with temperatures. It can be found from the composite configuration changes at failure that when the temperature is lower (100 K), the atomic structure is relative perfect at the beginning of tensile loading, the disordered atoms and subsequent micro-voids appear near the interface as the tensile process proceeds, the damage initiates from the micro-voids and propagates along the glide plane within the matrix, and finally, the Ni-CNT breaks near the interface pores. As the temperature increases from 100 to 500 K, there are more defects such as disordered atoms, point defects, dislocation, and slip bands in the Mg matrix and the near interface Mg matrix due to the increased atomic activity. In addition, the plastic deformation resistance for the Mg matrix and CNTs is weakened due to their structural changes during deformation at high temperatures [\[23](#page-10-22)]. The damage at high temperatures can be formed and propagated from multiple defects at the same time. As a consequent, the defective CNTs fracture occurs near the interface with more defects after undergoing a certain degree of necking.

3.3 Strain‑rate effect

Figure [9](#page-6-1) depicts the tensile stress–strain curves of Ni-CNT/ Mg nanocomposites at 300 K and three typical strain rates, in which the longitudinal coordinates represent the average

Fig. 9 Tensile stress–strain curves of Ni-CNT/Mg composites at 300 K and three typical strain rates

Fig. 8 Snapshots of tensile failure modes for Ni-CNT/Mg composites at different temperatures of **a** *T*=100 K, **b** *T*=200 K, **c** *T*=300 K, and **d** *T*=500 K

stress value of the Ni-CNT/Mg composites in the *z* (tensile) direction. It can be seen from Fig. [9](#page-6-1) that the stress–strain curves for the Ni-CNT/Mg composites at different strain rates are completely coincident in the initial linear elastic stage during their tensile loading, indicating no obvious change in Young's modulus of the composites with strain rate. This may be attributed to the little strain-rate sensitivity of the elastic modulus of nanomaterials or is slightly decreased with the strain rate increases [\[24](#page-10-23)]. In addition, it is worth noting that there all exists a secondary declined stress stage in the stress–strain curves of the composites at different strain rates after the peak stress; however, at the lower strain rate $(4 \times 10^8 \text{ s}^{-1})$, the fluctuations in the stress–strain curves become larger after the second stress drop due to the Ni-CNT fracture, while at the higher strain rate $(2 \times 10^{10} \text{ s}^{-1})$, the tensile stress decreases more slowly after it reaches the peak stress and the stress–strain curve is relatively smooth, so the peak stress and the corresponding strain is obviously increased. The variation in the stress–strain response after the elastic stage is related to the plastic deformation mechanism of the composites at different strain rates. The reasons for the different deformation mechanisms will be further analyzed in the following text.

The strain rate also influences the peak stress, corresponding tensile strain, fracture stress, and fracture strain of the Ni-CNT/Mg composites. Figure [10](#page-7-0) presents the variations in the peak stress, corresponding tensile strain, fracture stress, and fracture strain as a function of logarithmic strain rate for tensile deformation at 300 K. It can be seen from Fig. [10a](#page-7-0), b that when the strain rate is lower than 1×10^9 s⁻¹, the peak stress, corresponding tensile strain, and fracture strain have little strain-rate sensitivity and their values are kept at 6.95 GPa, 9.5 and 11.5%, respectively. On the other side, when the strain rate is higher than 1×10^9 s⁻¹, the peak stress and corresponding tensile strain increase continuously with the strain rate, showing an obvious strain-rate hardening effect. Meanwhile, the fracture strain also increases

as the strain rate increases, but the increment magnitude (56.67%) is larger than that (33.69%) of the tensile strain at point corresponding to the peak stress when the strain rate increases from 1×10^9 to 2×10^{10} s⁻¹. This may be related to the nonlinear stress–strain relation at the later stage of the tensile stress–strain curve, which means that a small change in the fracture stress will result in a big change in the fracture strain. In contrast to the peak stress and corresponding tensile strain as well as fracture strain, the fracture stress shows a fluctuating increase with the increasing strain rate over the entire strain-rate range (Fig. [10b](#page-7-0)). This increasing trend in fracture stress can be attributed to the combined effects of strain-rate hardening and more defects at high strain rates. From the above, it can be concluded that increasing strain rate has the effect of improving the peak stress and fracture stress, but the strain-rate effect on the mechanical properties is less significant compared to the temperature effect [\[25](#page-10-24)], as presented in the previous section.

3.4 Tensile response and fracture at different temperatures and strain rates

Figure [11](#page-8-0) shows the tensile stress–strain curves of Ni-CNT/ Mg nanocomposites under the coupled effects of different temperatures (200, 300, 400, and 500 K) and strain rates $(1 \times 10^9, 4 \times 10^9, 8 \times 10^9, \text{ and } 2 \times 10^{10} \text{ s}^{-1})$. It is noted that there is a general increase in tensile stress with increasing strain rate, but a decrease with increasing temperature, indicating that the composites have strain-rate hardening and temperature softening effects. However, the temperature has a greater effect on stress at low strain rate than that at high strain rate, for instance, when the temperature is increased from 200 to 500 K, the peak stress is decreased by 33.84% at a lower strain rate of 1×10^9 s⁻¹ but only 28.81% at a higher strain rate of 2×10^{10} s⁻¹. In addition, the stress–strain curves for the composites at different temperatures are a little different in shape with the increase of strain rate, especially the

Fig. 10 Correlations between mechanical properties and strain rates in Ni-CNT/Mg composites at 300 K: **a** peak stress and corresponding strain and **b** fracture stress and fracture strain as a function of logarithmic strain rate

Fig. 11 Tensile stress–strain curves of Ni-CNT/Mg composites at the temperatures and strain rates of **a** $T = 200-500$ K and $\epsilon = 1 \times 10^9$ s⁻¹, **b** *T*=200–500 K and ϵ = 4 × 10⁹ s⁻¹, **c** *T*=200–500 K and ϵ = 8 × 10⁹ s⁻¹, and **d** *T*=200–500 K and ϵ = 2 × 10¹⁰ s⁻¹

fracture stress and strain at low temperature show a general increasing trend with increasing strain rate.

Figure [12](#page-9-0) shows the snapshots of tensile failure modes for Ni-CNT/Mg composites at the fracture strains under the coupled effects of different temperatures and strain rates. It can be seen that the failure processes of the composites under the coupled effects of different temperatures and strain rates are approximately the same to some degree. These processes involve initial microvoid formation, matrix cracking and crack propagation along the microvoid, and final CNT fracture. However, the deformation and failure mechanisms of the composites at different temperatures and strain rates are different.

As shown in Fig. [12a](#page-9-0), the composite shows in general a "brittle fracture" mode and its fracture mainly occurs near the Ni-CNT/Mg interface at relatively high temperature and low strain rate. Under the above conditions, the disordered atoms and subsequent micro-voids preferentially appear near the interface, and the damage propagates along the glide plane inside the matrix and more dislocations are generated nearby. At last, the Ni-CNT fracture occurs at the expanded micro-voids near the interface due to the simultaneous bond breaking in it, and the composite has a lower failure strain. As shown in Fig. [12b](#page-9-0), c, when the temperature decreases and the strain rate increases, there are multiple fracture sites in the matrix and Ni-CNT as well as their interface during the failure of the composites, and the corresponding failure strain increases. The composite shows the "quasi-brittle fracture" characteristics. When the temperature is relative low and the strain rate is relative high, as shown in Fig. [10d](#page-7-0), the composite is characterized by a "mixed fracture". Under this condition, there are more slip bands, local amorphization and micro-voids in the matrix, and multiple damage failures have occurred in the matrix before the damage occurs at the interface. As a result, the multiple fracture of Ni-CNT occurs at the interface after undergoing a larger necking and the corresponding failure strain of the composite has the maximum value.

To further investigate the coupled effects of temperature and strain rate, the strain-rate sensitivity exponents at different temperatures are calculated and analyzed. The strainrate sensitivity exponent (*m*) measures how much a material changes in peak stress as strain-rate changes and it can be defined as follows [[25\]](#page-10-24):

Fig. 12 Snapshots of tensile failure modes for Ni-CNT/Mg composites under the coupled effects of different temperatures and strain rates: **a** $T = 500$ K and $\epsilon = 1 \times 10^9$ s⁻¹, **b** $T = 400$ K and $\epsilon = 4 \times 10^9$ s⁻¹, **c** $T = 300$ K and $\epsilon = 8 \times 10^9$ s⁻¹, and **d** $T = 200$ K and $\epsilon = 2 \times 10^{10}$ s⁻¹

Fig. 13 Strain-rate sensitivity exponents of the peak stress in the plot of $ln(\sigma)$ vs. $ln(\epsilon)$ for Ni-CNT/Mg composites at different temperatures. The slope of the fitting line is the strain-rate sensitivity exponent

 $m = \frac{\partial \ln(\sigma)}{\partial \ln(\dot{\epsilon})}$ (14)

Equation (14) (14) shows that the strain-rate sensitivity coefficient *m* can be obtained from the slope of the fitting line of $\ln(\sigma)$ vs. $\ln(\dot{\epsilon})$, as shown in Fig. [13.](#page-9-2) It can be seen that the strain-rate sensitivity exponent increases from 0.01678 to 0.04537 when the temperature ranges from 200 to 500 K. This indicates that the peak stress is more sensitive to strain rate at high temperatures. In the same way, it also can be found that the temperature effect on the peak stress is more significant at a low strain rate than that at a high strain rate. These conclusions are consistent with the summary results in Fig. [12.](#page-9-0)

4 Conclusions

- 1. MD simulations of uniaxial tension of Ni-CNT/Mg composites at different temperatures and strain rates have shown that the Ni-coated CNT can effectively enhance the elastic modulus and peak stress of the composites.
- 2. Temperature and strain rate have a strong influence on the tensile mechanical properties of Ni-CNT/Mg composites. The peak stress decreases by about 42.99% at 1×10^{9} s⁻¹ when the temperature increases from 100 to

500 K, while the peak stress of Ni-CNT/Mg composite at the room temperature (300 K) increases by 12.76% as the strain rate increases from 1×10^9 to 1×10^{10} s⁻¹. Moreover, the mechanical properties of the composites are affected by the coupled effect of temperature and strain rate. The mechanical properties are more sensitive to high temperature at low strain rate, and reciprocally, to high strain rate at low temperature.

3. The Ni-CNT/Mg composite shows different failure modes under the coupled effects of different temperatures and strain rates. The composites show good overall mechanical behavior at low temperature and high strain rate.

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References

- 1. A. Dey, K.M. Pandey, Magnesium metal matrix composites—a review. Rev. Adv. Mater. Sci. **42**, 58–67 (2015)
- 2. T.S. Srivatsan, C. Godbole, T. Quick, M. Paramsothy, M. Gupta, Mechanical behavior of a magnesium alloy nanocomposite under conditions of static tension and dynamic fatigue. J. Mater. Eng. Perform. **22**, 439–453 (2013)
- 3. N. Silvestre, State-of-the-art review on carbon nanotube reinforced metal matrix composites. Intl. J. Compos. Mater. **3**, 28–44 (2013)
- 4. J. Liang, H. Li, X. Hu, J. Wei, L. Qi, Fabrication of Ni-coated carbon nanotubes reinforced magnesium matrix composites, In: International Conference on Manipulation, Manufacturing and Measurement on the Nanoscale, IEEE, 2014, 299–302
- 5. M.H. Nai, J. Wei, M. Gupta, Interface tailoring to enhance mechanical properties of carbon nanotube reinforced magnesium composites. Mater. Design. **60**, 490–495 (2014)
- 6. X. Zhou, S.Y. Song, L. Li, R.J. Zhang, Molecular dynamics simulation for mechanical properties of magnesium matrix composites reinforced with nickel-coated single-walled carbon nanotubes. J. Compos. Mater. **50**, 191–200 (2016)
- 7. H.Y. Song, X.W. Zha, Mechanical properties of nickel-coated single-walled carbon nanotubes and their embedded gold matrix composites. Phys. Lett. A. **374**, 1068–1072 (2010)
- 8. K. Duan, L. Li, Y.J. Hu, X.L. Wang, Damping characteristic of Ni-coated carbon nanotube/copper composite. Mater. Design. **133**, 455–463 (2017)
- 9. LAMMPS Molecular Dynamics Simulator. Available at: [http://](http://www.cdc.gov/kawasaki/) www.cdc.gov/kawasaki/. Accessed January 9, 2018
- 10. S. Inoue, Y. Matsumura, Molecular dynamics simulation of physical vapor deposition of metals onto a vertically aligned

single-walled carbon nanotube surface. Carbon **46**, 2046–2052 (2008)

- 11. D.Y. Sun, M.I. Mendelev, C.A. Becker, K. Kudin, T. Haxhimali, M. Asta, J.J. Hoyt, A. Karma, D.J. Srolovitz, Crystal-melt interfacial free energies in hcp metals: A molecular dynamics study of Mg. Phys. Rev. B **73**, (2006). 024116-1-024116-12
- 12. S.J. Stuart, A.B. Tutein, J.A. Harrison, A reactive potential for hydrocarbons with intermolecular interactions. J. Chem. Phys. **112**, 6472–6486 (2000)
- 13. M. Sammalkorpi, A. Krasheninnikov, A. Kuronen, K. Nordlund, K. Kaski, Mechanical properties of carbon nanotubes with vacancies and related defects. Phys. Rev. B **70**, 121–127 (2004)
- 14. Y. Shibuta, S. Maruyama, Bond-order Potential for transition metal carbide cluster for the growth simulation of a single-walled carbon nanotube. Comp. Mater. Sci. **39**, 842–848 (2007)
- 15. A. Stukowski, Visualization and analysis of atomistic simulation data with OVITO-the Open Visualization Tool. Model. Simul. Mater. Sci. Eng. **18**, 015012 (2010)
- 16. S. Lee, J. Park, J. Yang, W. Lu, Molecular dynamics simulations of the traction-separation response at the Interface between PVDF binder and graphite in the electrode of Li-ion batteries. J. Electrochem. Soc. **161**, A1218–A1223 (2014)
- 17. L. Li, M. Han, Molecular dynamics simulations on tensile behaviors of single-crystal bcc Fe nanowire: effects of strain rates and thermal environment. Appl. Phys. A **123**, 450 (2017)
- 18. J.F. Xiang, L.J. Xie, S.A. Meguid, S.Q. Pang, J. Yi, Y. Zhang, R. Liang, An atomic-level understanding of the strengthening mechanism of aluminum matrix composites reinforced by aligned carbon nanotubes. Comp. Mater. Sci. **128**, 359–372 (2017)
- 19. F.L. Zhu, H.Y. Liao, K. Tang, Y.K. Chen, S. Liu, Molecular dynamics study on the effect of temperature on the tensile properties of single-walled carbon nanotubes with a Ni-coating. J. Nanomater. 1–7 (2015) (2015)
- 20. N. Silvestre, B. Faria, J.N.C. Lopes, Compressive behavior of CNT-reinforced aluminum composites using molecular dynamics. Comp. Sci. Technol. **90**, 16–24 (2014)
- 21. A.L. Kalamkarov, A.V. Georgiades, S.K. Rokkam, V.P. Veedu, M.N. Ghasemi-Nejhad, Analytical and numerical techniques to predict carbon nanotubes properties. Int. J. Solids. Struct. **43**, 6832–6854 (2006)
- 22. L. Liao, Q.H. Yuan, Microstructure and mechanical properties of AZ91 alloy composites reinforced with Ni-coated carbon nanotubes. Mater. Mech. Eng. **34**, 92–95 (in Chinese) (2010)
- 23. S.F. Hassan, M. Paramsothy, Z.M. Gasem, F. Patel, M. Gupta, Effect of carbon nanotube on high-temperature formability of AZ31 magnesium alloy. J. Mater. Eng. Perform. **23**, 2984–2991 (2014)
- 24. Y.G. Zhou, M. Hu, Mechanical behaviors of nanocrystalline Cu/ SiC composites: An atomistic investigation. Comp. Mater. Sci. **129**, 129–136 (2017)
- 25. M.Q. Chen, S.S. Quek, Z.D. Sha, C.H. Chiu, Q.X. Pei, Y.W. Zhang, Effects of grain size, temperature and strain rate on the mechanical properties of polycrystalline graphene – A molecular dynamics study. Carbon **85**, 135–146 (2015)